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M. El-Kady

Mechanical Power Engineering Department., Faculty of Engineering., El-Mansoura University., Mansoura., Egypt, mselkady@mum.mans.eun.eg

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# Natural Convection From Dual Surface Heat Sources in a Vertical Rectangular Enclosure

الحمل الطبيعي من منبعين حراريين سطحيين في حيز رأسي مستطيل

M.S. El-Kady Mechanical Power Engineering Department Mansoura University, Egypt

#### خلاصـــــة:

#### Abstract:

Natural convection in a heated two-dimensional rectangular vertical enclosure is investigated numerically and experimentally for different dual heater configurations. The governing conservation equations of mass, momentum, and energy for the problem are solved numerically using the finite difference technique. The study is carried out for the steady laminar flow of air in an enclosure with constant heat flux dual heaters at one adiabatic insulated vertical wall, an isothermally cooled other vertical wall and insulated horizontal walls. The range of parameters considered is  $0.1 \le S_1/H \le 0.5$ ,  $0.1 \le S_2/H_1 \le 0.875$ ,  $1 \le A \le 10$ .  $10^4 \le Ra \le 10^7$  and Pr = 0.7. The results show that for the lower heater  $\theta_{max}$  and Nu are not affected by the location of the upper heater and the aspect ratio for  $A \ge 3$ . For the upper heater,  $\theta_{max}$  and Nu are independent on the aspect ratio and the upper heater location in the optimum range of the upper heater location  $0.1 \le S_1/H_{\le} 0.5$ , the upper heater location  $S_2/H_1 \le 0.5$  and  $A \ge 3$ .

for bottom heater:  $\theta_{max} = 1.055 \text{ Ra}^{-0.139}$ ,  $\overline{Nu} = 1.4 \text{ Ra}^{0.1179}$ for top heater:  $\theta_{max} = 1.4 \text{ Ra}^{-0.153}$ ,  $\overline{Nu} = 0.99 \text{ Ra}^{0.135}$ Experimental tests were done for A=2, S<sub>1</sub>/H= 0.2, S<sub>2</sub>/H<sub>1</sub>= 0.29, W/L=2.5 and Ra=3.2x10<sup>4</sup>, 1.23x10<sup>5</sup>, and 2.65x10<sup>5</sup> and show the validity of the numerical simulation.

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M. 41

#### Introduction

In recent years, the study of natural convection in enclosures has become increasingly important. In particular, within the development of complex, high power electronics packaging and with the increasing demands on exhausting energy supplied, the understanding of free convective flows within rectangular enclosures heated by concentrated discrete heat sources has become warranted. In modern electronic equipments, a large number of high power dissipating components such as transistors, resistors, and power transformers are being packed in modular rectangular enclosures. The applications of these packages are often such that space, weight, and external cooling sources are minimal. In these cases, the placement of the high power dissipating components within an electronic package should be optimized as to maximize the natural convection heat transfer within the enclosure, thus possibly totally eliminating the need for forced cooling.

Extensive surveys of the various modes of convective heat transfer and relevant configurations along with the associated heat transfer and other correlations have been presented by Jaluria [1], Incropera [2], Papanicolaon and Jaluria [3] and El Kady and Araid [4]. Several works in the literature are dealing with mixed convection in the cooling of protruding heat sources of electronic components. among them the work of Habchi and Archarya [5], Kang et. al. [6], Mahaney et al. [7], Kim et. al. [8], and Papanicolaou and Jaluria [3]. Forced convection has been studied by many authors. Among them were Incropera et al. [9], Davalath and Bayazitoglu [10] and Kim and Anand [11-13], and Nakayama and Park [14]. Natural convection cooling techniques have distinct advantages because of their low cost, ease of maintenance and absence of electromagnetic interference and operating noise. Jaluria studied the buoyancy-induced flow due to isolated thermal sources on a vertical plate [15]. The natural convection hetween series of vertical plate channels with embedded line heat sources has been studied in order to meet the relatively lower operating temperature requirements of integrated circuits by Anand et al. [16], and Kim et al. [17,18]. Another important configuration involve complete and partial enclosures. Such as those that are encountered in small electronic devices and personal computers. Not much work has heen done on such enclosure flows. But there is a growing interest in these problems. Keyhani ct al. [19] investigated experimentally natural convection flow and heat transfer characteristics of an array of discrete heat sources in enclosures, and in [20] they showed the aspect ratio effect in natural convection. Carmona and Keyhani [21] showed the cavity width effect on cooling of five flush heaters on one vertical wall of an enclosure. El Kady and Araid [4] showed the effect of size and location of a surface heater embedded in the vertical wall of a two dimensional rectangular enclosure.

It is important to determine the extent of the wake region above the isolated heat sources, so that other components and surface heaters may be located without substantially reducing the heat transfer from them. The present study is directed to develop numerical model to investigate the heat transfer characteridtics due to the natural convection from two-dimensional discrete heat sources embedded in the vertical wall of the rectangular enclosures. It concentrates on the effect of the locations of the two heat sources, the eoclosure aspect ratio and Rayleigh number on the surface temperature of the heat sources and the local and average heat transfer from them. Experimental test was done and the results were used to validate the numerical model.

#### Mathematical Formulation and Numerical Procedure

Two-dimensional steady laminar natural convection air flow is considered inside a vertical rectangular enclosure of height H and width W, as illustrated in Fig. 1. Two flat heat sources, with uniform heat flux input q and height L, are located on vertical adiabatic surface. The equations describing the problem are the Navier-Stokes equations for the fluid, with buoyancy effects taken into account and using the Boussinesq approximations, as well as the the energy equation, which describes the temperature variation through the enclosure (fluid and solid walls). With the introduction of the stream function and the vorticity as the independent variables, the nondimensional equations can be written as :

Vonicity Eq.:	U.∂ω/∂X + V	$\partial \omega / \partial Y =$	$\partial^2 \omega / \partial X^2 + \partial^2 \omega / \partial Y^2 - $	· R₄/Pr.[∂θ/∂X]	(1)
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 $-\omega = \frac{\partial^2 \Psi}{\partial X^2} + \frac{\partial^2 \Psi}{\partial Y^2}$ Stream function equation: (2)

Energy equation: 
$$U.\partial\theta/\partial X + V.\partial\theta/\partial Y = 1/\Pr[\partial^2\theta/\partial X^2 + \partial^2\theta/\partial Y^2]$$
 (3)

where 
$$U = \partial \Psi / \partial Y$$
 and  $V = -\partial \Psi / \partial X$  (4)

 $\omega$  is the dimensionless vorticity ( $\omega = \partial V / \partial X \cdot \partial U / \partial Y$ ),  $\Psi$ ,  $\theta$ , U, V, and v are the dimensionless stream function, dimensionless temperature, dimensionless velocity in the dimensionless horizontal and vertical directions X and Y, and the fluid kinematic viscosity, respectively.

In the nondimensionalization, the height of the heater L, is taken as the characteristic length scale. The velocity scale is the  $L/\nu$ , and the temperature scale is  $\Delta T = q L/k$ . The dimensionless parameters are the Prandtl number  $Pr = v/\alpha$  and the Rayleigh number  $Ra = g\beta L^4 q/(k\alpha v)$ . Where k,  $\alpha$ , and  $\beta$  are the thermal conductivity, thermal diffusivity, and thermal expansion coefficient, respectively. More details for the formulation can be found in El Kady and Araid [4].

The dimensionless boundary conditions for the present system are

$$U = V = \Psi = 0$$
, on all boundaries (5a)

at X= 0: 
$$\partial \theta / \partial X = -1$$
 at heater surface, and  $\partial \theta / \partial X = 0$  elsewhere (5b)  
 $\theta = 0$  at X= W/L, and  $\partial \theta / \partial Y = 0$  at Y=0 or H/L. (5c)

$$\Theta = 0$$
 at  $X = W/L$ , and  $OO/OT = 0$  at  $T = 0.01$  F/L (5)

The local and average heat transfer coefficients along the surface of the heater are presented by means of Nusselt number and average Nusselt number according to its definition for isoflux heating as:

$$Nu = h.L/k = (q.L/k)/(T-T_w)$$
(6)

$$Nn = (q L/k)/(T - T_w)]$$
(7)

where Tw and T are the cooling wall temperature and the average heater surface temperature, respectively. Nu and Nu are the reciprocal of the dimensionless surface local and average temperatures, respectively.

The dimensionless governing equations (1)-(4) and the associated boundary conditions given by equation (5) were solved by a finite difference procedure discussed by Patanker [22]. Both the first and second order derivatives were discretized by using eentral difference formulas except the convection terms which were discretized by using the upwind scheme. The difference algebraic equations are solved using the Gauss-elimination method. An iterative solution procedure was employed here to obtain the steady state solution of the problem considered. A non-uniform grid with

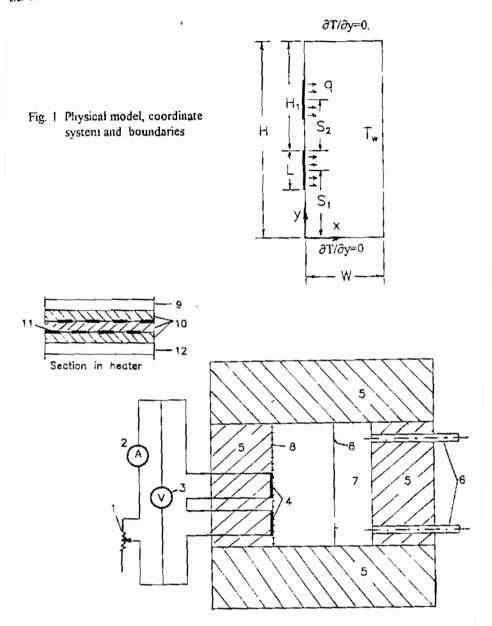


Fig. 2 Schematic diagram of the experimental apparatus

(1) auto transformer, (2) ammeter, (3) voltmeter, (4) electric heaters, (5) Polystyrene insulation, (6) inlet and outlet cooling water tubes, (7) Heat exchanger, (8) copperconstantan thermocouple, (9) stainless sheet plate, (10) sheet of mica, (11) nickelchromium electric heater, (12) Thermal insulation layer denser clustering near the walls was considered to give grid independent results. It varies from 41x41 for A=1 to 21x101 for A=10. The convergence criteria is attained when the relative error between two successive iterations for the field variables  $\omega$ ,  $\Psi$ , or  $\theta$  is less than  $\xi = 10^{-5}$ 

#### Experimental Work:

To validate the program developed in this work an experimental study for the air flow in a rectangular enclosure is carried out. A schematic diagram of the designed and constructed apparatus is shown in Fig. 2. The used two dimensional rectangular enclosure is 100 mm height and 50 mm width and 300 mm in the spanwise direction. The right vertical wall of the enclosure was held at constant temperature by fitting a copper counter-flow heat exchanger (7) in which the city water was used as a coolant liquid. The two horizontal surfaces were maintained adiabatic by conducting them of closed-pore extruded polystyrene insulation with 50 mm thickness. The left vertical wall of the enclosure was made of the closed-pore extruded polystyrene insulation (5) of 50 mm thickness. Two electric heaters (4) with dimensions of 20 mm height, 2.5 mm thickness, and 300 mm long was embedded in the left vertical adiabatic wall to make the heater face in the same vertical inside plane of the wall. The vertical distances between the two heaters was 10 mm, between the bottom horizontal wall and the lower heater was 10 mm and between the upper heater and the top horizontal wall was 40 mm. Each heater face is made of polished stainless steel sheet (9) with 0.5 mm thickness. This sheet was heated electrically by an electric heater. A nickel-chromium wire (11) was wounded around a mica sheet (10) of 0.5 mm thickness and then was sandwiched hetween other two mica sheets (10). The heat input to the heaters was controlled by using an auro-transformer (1), an animeter (2) and voltmeter (3).

The temperature distribution along the heated and unheated parts of the vertical wall was measured by 24 copper-constantan thermocouples. 20 thermocouples were fixed in the middle plane of the spanwise as shown in Fig. 2. The other 4 thermocouples were fixed in the middle height of each stainless sheet at 100 nm distance in both sides from the middle plane of the spanwise. Two thermocouples (8) were embedded in the wall of the heat exchanger (7) and monitored to ensure a uniform temperature distribution. The thermocouples were connected to digital temperature recorder. Nearly three hours were needed to reach the steady state condition which was recorded, as the temperature reading did not change within a time of about 15 minutes.

#### **Results and Discussion**

A careful examination of the model equations and the associated houndary conditions reveals that the independent parameters are the fluid Prandtl number Pr, the Rayleigh number Ra, the aspect ratio A=H/W, the location of the lower heater  $S_1/H$ , and the wake distance between the upper and lower heater,  $S_2/H_1$ . El Kady and Araid [4] concluded that the optimum location for a single heater in a vertical enclosure corresponding to maximum heat transfer and minimum surface temperature lies in the

M.45

# M. 46 M.S. El-Kaily

range of  $0.1 \le S_1/H \le 0.5$ , i.e. i) the lower half of the enclosure. Therefore, the location of the lower heater is chosen to be in this optimum range in the present study. For the  $a^{1/3}$  of Pr=0.7, the present study concentrates on the effect of the variation of the upper heater location (wake distance between upper and lower heaters)  $0.125 \le S_2/H_1 \le 0.875$ , the variation of the aspect ratio  $1 \le A \le 10$ , and the variation of Rayleigh number  $10^4 \le Ra \le 10^7$  on the flow and heat transfer characteristics to get the optimum position for the upper heater corresponding to the maximum heat transfer and minimum heater surface temperatures and correlations s = - describes these variables.

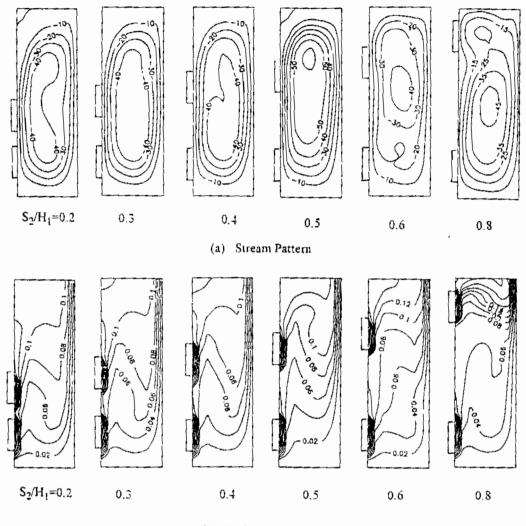
#### Effect of the upper scatter location

The heat transfer from the heaters and the surface temperatures of them depend mainly on the short, and temperature of the air flowing in the adjacent boundary layer. So, the m - single of the behaviour of the streamlines and the isothermal contours inside the enclosure is necessary

Figure 3 present the streamlines and temperature contours for  $S_1/H = 0.1667$ , Ra = 10<sup>5</sup> with the section of the upper heater location relative to the lower heater  $S_2/H_1 = 0.2$ , 0.3 + 0.0 and 0.8. Near the heated section along the hot wall and the opposing vertical state mick layers of primary flows are observed in Fig. 3a. Between these layers expected and constructions from the low region of the lower heater and enors near the top surface of the enclosure. For  $S_2/H_1 = 0.8$  the upper heater moves nearer to the top of the enclosure, secondary cell in front of the upper heater exists, which in turns causes a decrease in the maximum stream function and the flow velocity in front of the upper heater. As shown in Fig. 3b, a growth of the thermal boundary layer along the two heaters exists. The temperature distribution for the core fluid and the cold wall is nearly constant with the exception of the top section of the enclosure. In the top section of the enclosure the temperature increases as the upper heater moves towards the top horizontal adiabatic surface. The temperature in front of the lower acater  $\sim$  hearing constant and independent on the location of the upper heater.

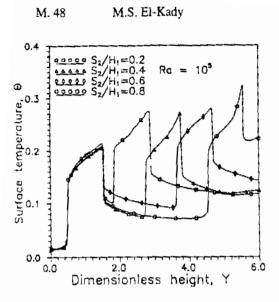
Fig. 4 shows the temperature distribution at the hot wall for the heated and unheated sections, while  $\log |s|_{presents}$  the local Nusselt number for both heaters for  $S_1/H=0.1667$ ,  $R_0=0.05$  and d=0.05, 0.4, 0.6 and 0.8. The above characteristics of the steamlines and usorber denoted the enclosure reveals that the air coming from the cooled wall and exemption is detailing edge of the lower heater is nearly at constant temperatures locall constant d bases. Therefore, the surface temperature and the local Nu of the lower heater is detailed or the four cases, the surface temperature of the lower heater is maller than the or the upper heater, and the local Nusselt number for the lower heater is high a than that of the upper heater for the four cases considered in Figs. 4 and 5

The influence of the upper heater location (wake distance)  $0.125 \le S_2/H_1 \le 0.875$ on the maximum surface temperatures and the average Nusselt number for both heaters are presented in Figs mand? The maximum surface temperature  $\theta_{max}$  for the lower heater is smaller than that of the upper heater while Nu is higher than that of the upper heater. For the lower heater both the maximum surface temperature  $\theta_{max}$  and the average Nusselt number. No are not affected by the location of the upper heater (the variation of  $S_2/H_1$ ). For the upper heater, with the increase of  $S_2/H_1$  (the wake



(b) Isotherms Contours

Fig. 3 Stream pattern and isotherm contours for different upper heater locations, Pr = 0.7,  $S_1/H=0.1667$ . A=3, and Ra=10<sup>5</sup>



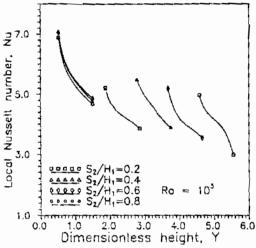


Fig. 4 Surface temperature distribution for different upper heater locations,  $P_T = 0.7$ ,  $S_1/H=0.1667$ , A=3, and  $Ra=10^5$ 

Fig. 5 Local Nusselt Number variation for different upper heater locations, Pr = 0.7,  $S_1/H=0.1667$ , A=3, and  $Ra=10^5$ 

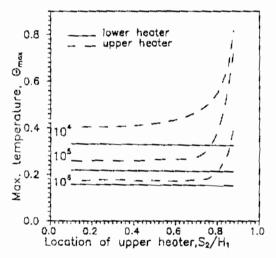


Fig. 6 Influence of upper heater locations on the maximum surface temperature for Pr = 0.7,  $S_1/H=0.1667$ , A=3

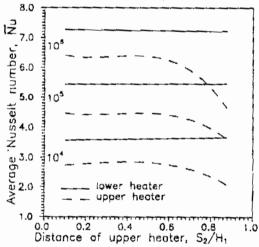


Fig. 7 Influence of the upper heater locations on the average Nusselt number for Pr=0.7,  $S_1/H=0.1667$ , A=3

distance increases), both  $\theta_{max}$  and Nu are independent on the upper heater location and are nearly constant until  $S_2/H_1 = 0.6$ . For  $S_2/H_1 \ge 0.6$ , the upper heater becomes nearer to the top surface, the maximum stream function in front of it decreases and the air flow passing through the adjacent boundary layer becomes slower with increasing temperature. For those reasons, the maximum surface temperature  $\theta_{max}$  increases sharply and the average Nusselt number Nu sharply decreases. Therefore, the optimum location for the upper heater corresponding to maximum average heat transfer Nu and minimum surface maximum temperature  $\theta_{max}$  lies in the lower half of the distance above the trailing edge of the lower heater where  $S_2/H_1 \le 0.6$ .

# Effect of Aspect Ratio

Figure 8 presents the streamlines and isothermal contours with the variation of aspect ratio  $2 \le A \le 8$  for Ra = 10<sup>5</sup>, while Figs. 9 and 10 present the variation of the maximum surface temperature and average Nusselt number for both the lower and upper heaters with the variation of the aspect ratio  $1 \le A \le 10$  for Ra =  $10^4$ ,  $10^5$ , and  $10^6$ . In Figs 8-10 both heaters lie in the optimum location range with  $S_1/H= 0.1667$  and  $S_2/H_1 = 0.3$  As shown in Fig. 8a, thick layers of primary flow are observed upar the hot and opposing vertical walls. With the increase of the aspect ratio a long and narrow core flow exists between the primary flow with disappearing of the secondary core flow eells. A secondary recirculating zone with low intensity can be predicted in the upper left corner indicating the separation of the flow originating from the top heater and rising along the remainder of the adiabatic wall. Although higher maximum streamfunctions appear with the increase of aspect ratio, the maximum values of streamfunctions are nearly constant for A23 in front of each heater which means nearly constant flow velocity inside the enclosure in front of each heater. Fig. 8h shows that with the increase of the aspect ratio, the active height of the isothermal cold wall increases and the chance to transfer more beat from the heater to the cold wall increases. By the Aspect ratio A≥3, the temperature of the core flow in front of each heater and temperature in the boundary layer adjacent to them are nearly constant. For the above reasons the air adjacent to the heaters flows with nearly constant velocity and temperature, and both the average Nusselt number and the maximum surface temperatures for both the lower and upper heaters are nearly constant and independent on the Aspect ratio for A≥3 as shown in Figs. 9 and 10. The same behaviour of the average Nusselt number and maximum surface heater temperatures exist for different values of Rayleigh number:  $Ra = 10^4$ ,  $10^5$ , and  $10^6$ .

#### Effect of Rayleigh number

For constant Rayleigh number, Figs. 7 and 8 show that, the maximum surface temperature and the average Nusselt number of both the lower and upper heaters are independent on  $S_2/H_1$  for the range of  $S_2/H_1 \le 0.6$ , while Figs. 9 and 10 present also that the maximum surface temperature and average Nusselt number of both the lower and upper heaters are independent on enclosure aspect ratio for the range of  $A \ge 3$ .

For the optimum range of the lower heater location  $0.1 \le S_1/H \le 0.5$ , the optimum range of upper location  $S_2/H_1 \le 0.6$  and the optimum range of enclosure aspect

M. 49

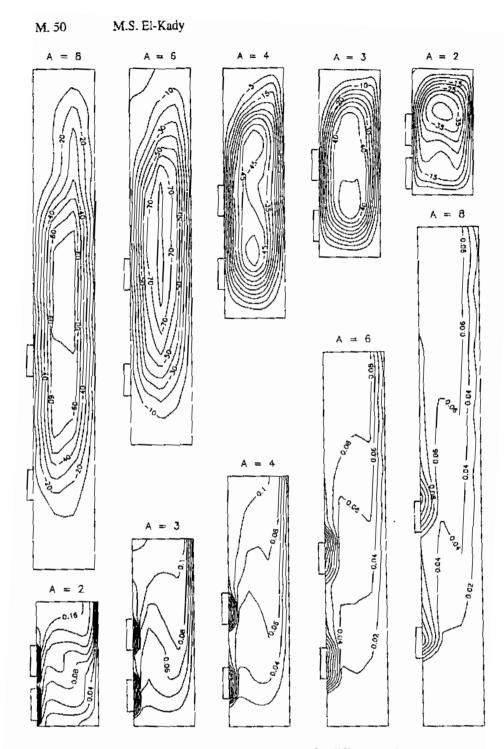


Fig. 8 Stream pattern and isotherm contours for different enclosure aspect ratios, Pr =0.7,  $S_1/H$ =0.1667,  $S_2/H_1$ =0.3, and Ra=10<sup>5</sup>

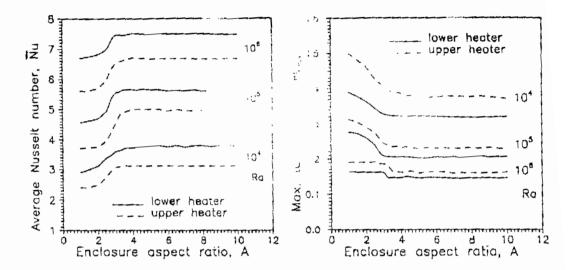


Fig. 9 Influence of enclosure aspect ratio on the average Nusselt number for Pr =0.7, S<sub>1</sub>/H=0.1667, S<sub>2</sub>/H<sub>1</sub>=0.3

Fig. 10 Influence of enclosure aspect ratio on the maximum surface temperature for  $P_T =$ 0.7, S<sub>1</sub>/H=0.1667, S<sub>2</sub>/H<sub>1</sub>=0.3

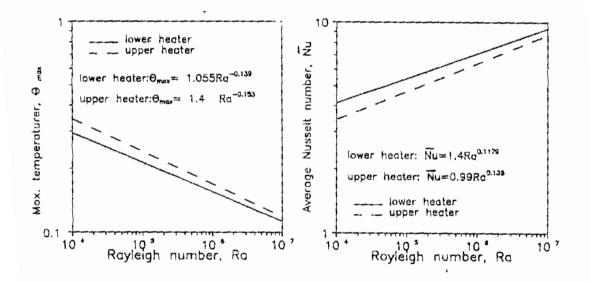


Fig. 11 Influence of Rayleigh number on the maximum surface temperature for  $Pr = 0.7, 0.1 \le S_1/H \le 0.5, S_2/H_1 \le 0.6$ , and  $A \ge 3$ 

Fig. 12 Influence of Rayleigh number on the average Nusselt number for Pr = 0.7,  $0.1 \le S_1/H \le 0.5$ ,  $S_2/H_1 \le 0.6$ , and  $A \ge 3$ 

ratio A≥3, Figs.11 and 12 present the variation of the maximum surface temperature and the average Nusselt number of both the bottom and top heaters with tha variation of Rayleigh number  $10^4 \le Ra \le 10^7$ .

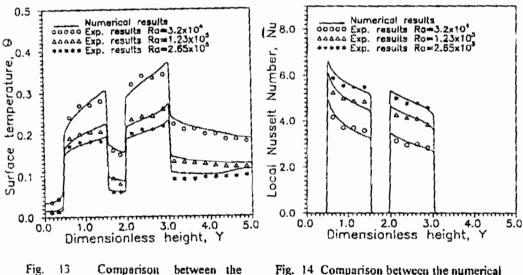
Figure 11 shows that the maximum surface temperature of the upper heater is always higher than that of the lower heater. Therefore, the average Nusselt number for the top heater is always smaller than that of the bootom heater. The maximum heater surface temperature and the average Nusselt number for both the upper and lower heaters are correlated as a function of Rayleigh number by equations as follows:

for bottom heater:	$0_{\rm max} = 1.055 \ {\rm Ra}^{-0.139}$	(8)
	$\bar{N}u = 1.4  Ra^{0.1179}$	(9)
for too heater:	$\theta_{\rm max} = 1.4  \mathrm{R}_{2} \cdot 0.153$	(10)

by top heater: 
$$\theta_{\text{inax}} = 1.4 \text{ Ra}^{-0.153}$$
 (10)

$$Nu = 0.99 Ra^{0.135}$$
(11)

To validate the numerical calculated results, experimental tests were done for A=2, W/L=2.5,  $S_1/H= 0.2$ ,  $S_2/H_1= 0.29$ , Pr=0.7 and Ra=3.2x10<sup>4</sup>, 1.23x10<sup>5</sup>, and 2.65x10<sup>5</sup>. Comparison is shown in Fig. 13 between the calculated local surface heated and unheated parts of the wall and the experimental results. Fig. 14 presents also a comparison between the local heater surface Nusselt number with the experimental results calculated from the experimental data according to Eq. 6. Good agreement exists between the predicted and experimental results. The average and maximum error between the predicted and experimental data are 5 and nearly 10 percent, respectively.



numerical and experimental surface temperature for A=2, W/L=2.5, Pr=0.7, S1/H=0.2, S2/H1=0.29

Fig. 14 Comparison between the numerical and experimental heater local Nusselt number for A=2, W/L=2.5, Pr=0.7, S<sub>1</sub>/H=0.2, S<sub>2</sub>/H<sub>1</sub>=0.29

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# Conclusions

The effect of the wake distance between upper and lower heaters, the enclosure aspect ratio, and Rayleigh number on the flow, heaters surface temperature and heat transfer characteristics are studied numerically. The range of parameters considered is Pr=0.7, lower heater location  $0.1 \le S_1/H \le 0.5$ , upper heater location  $0.1 \le S_2/H_1 \le 0.875$ , the aspect ratio  $1 \le A \le 10$ , and  $10^4 \le Ra \le 10^7$ . The range of parameters for the experimental study is Pr=0.7, A=2,  $S_1/H=0.2$ ,  $S_2/H_1=0.29$ , W/L=2.5 and  $Ra=3.2\times 10^4$ ,  $1.23\times 10^5$ , and  $2.65\times 10^5$ . The results show the following conclusions:

The average and maximum surface temperatures of the upper heater are always higher than those of the lower heater. Therefore, the local and average Nusselt number for the upper heater are always smaller than those of the lower heater.

For the lower heater both the maximum surface temperature  $\theta_{max}$  and the average Nusselt number  $\overline{N}u$  are not affected by the location of the upper heater and the aspect ratio for A $\geq 3$ .

For the upper heater,  $\theta_{max}$  and Nu are independent on the aspect ratio for A≥3, and the upper heater location for S<sub>2</sub>/H<sub>1</sub>≤0.6 and the optimum location for the upper heater lies in the lower half of the distance above the trailing edge of the lower heater where S<sub>2</sub>/H<sub>1</sub>≤0.6 and A≥3.

For the optimum ranges of the lower heater location  $0.1 \le S_1/H \le 0.5$ , the upper heater location  $S_2/H_1 \le 0.6$  and the enclosure aspect ratio  $A \ge 3$ , the maximum heater surface temperature and the average Nusselt number for both the upper and lower heaters are correlated with Rayleigh number as follows:

for bottom heater:	$\theta_{\rm max} = 1.055 \ {\rm Ra}^{-0.139},$	Nu =	I.4 Ra <sup>0 [179</sup>
for top heater:	$\theta_{\rm max} = 1.4 \ {\rm Ra}^{-0.153}$	Nu =	0.99 Ra <sup>0.135</sup>

#### Nomenciature

А	aspect ratio = $H/W$
â	gravitational acceleration, n/s <sup>2</sup>
Ĥ	euclosure height, in
H	distance above the lower heater, m
k.	thermal conductivity, W/mK
L	heater height, ut
Nu	local Nusselt number along the heater surface, Eq. (9)
Nu Nu	average Nusselt number, Eq. (10)
Р	pressure, Pa
Pr	Prandile number, $v/\alpha$
q	heat flux at the heater surface, $W/m^2$
Ra	Rayleigh number, $Ra = g \beta L^4 q / (k\alpha v)$
S	the height of the lower heater cemter, m
S <sub>1</sub>	height of the upper heater eenter above the lower heater trailing edge, m
т	temperature, K

#### M. 54 M.S. El-Kady

u. v	field velocities in the x and y directions, m/s
U, V	non-dimensional field velocities in the X and Y direction, $uL/v$ , $vL/v$
W	enclosure width, m
х, у	distance in horizontal and vertical directions, m
X, Y	dimensionless distance in x direction, $x/L$ and $y/L$
α	thermal diffusivity eoefficient, m <sup>2</sup> /s
β	volumetric coefficient of thermal expansion, 1/K
θ	dimensionless temperature, $(T-T_w)/(q.L/k)$
Ψ	stream function, m <sup>2</sup> /s
Ψ	nondimensional stream function, $\psi/v$
Ç	vorticity, 1/s
ω	nondimensional vorticity, CL <sup>2</sup> /v
v	kinematic viscosity of the fluid, m <sup>2</sup> /s
ρ	fluid density, kg/m <sup>3</sup>

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